

dc SQUID

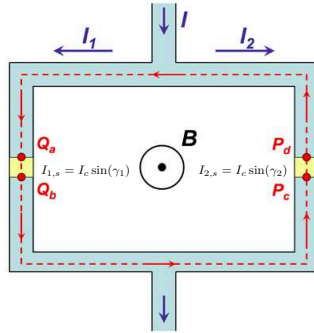
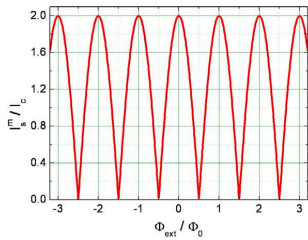
Negligible screening, zero voltage

$$\beta_L = \frac{2LI_c}{\Phi_0} \ll 1$$

Circulating currents do not compensate the external flux

$$\Phi \simeq \Phi_a$$

$$I_{s,max} \simeq 2I_c \left| \cos \left(\pi \frac{\Phi_a}{\Phi_0} \right) \right|$$



Figures from:
R. Gross, A. Marx, Walther Meissner Institut
<http://www.wmi.badw.de/teaching/LectureNotes/>

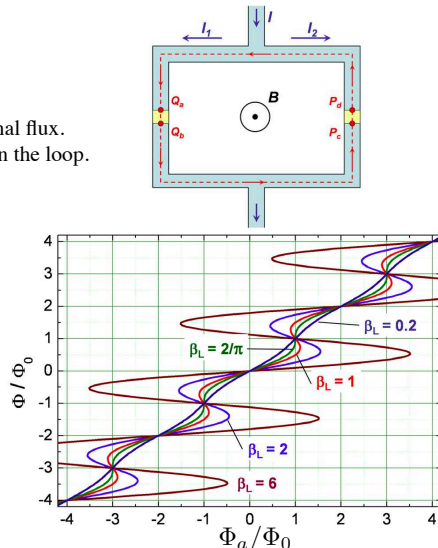
Nonzero screening, zero voltage

$$\beta_L = \frac{2LI_c}{\Phi_0}$$

Circulating currents compensate the external flux.
Strong screening: towards flux quantization in the loop.

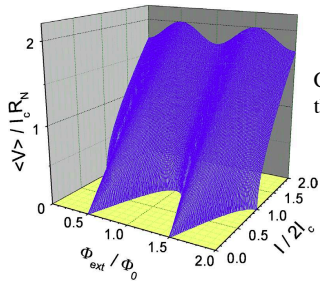
$$\Phi \neq \Phi_a$$

$$\frac{\Phi_a}{\Phi_0} = \frac{\Phi}{\Phi_0} + \frac{\beta_L}{2} + \sin \left(\pi \frac{\Phi}{\Phi_0} \right)$$



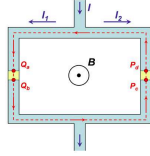
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Negligible screening, nonzero voltage

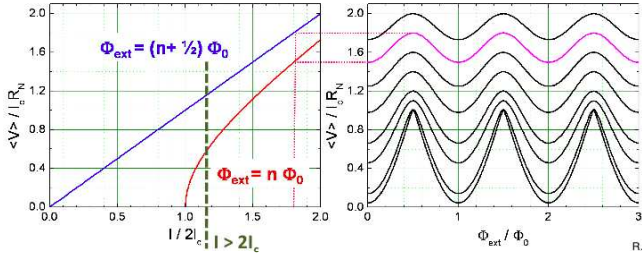


$$\beta_L = \frac{2LI_c}{\Phi_0} \ll 1$$

Circulating currents do not compensate the external flux $\Phi \simeq \Phi_a$



$$\langle V(t) \rangle = I_c R_N \sqrt{\left(\frac{I}{2I_c}\right)^2 - \cos^2\left(\pi \frac{\Phi_a}{\Phi_0}\right)}$$



Maximum modulation of the average voltage for $I = 2I_c$.

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dc SQUID: transfer function

$$\langle V(t) \rangle = I_c R_N \sqrt{\left(\frac{I}{2I_c}\right)^2 - \cos^2\left(\pi \frac{\Phi_a}{\Phi_0}\right)}$$
 Flux-to-voltage converter.

Transfer function: $H = \left| \frac{\partial \langle V(t) \rangle}{\partial \Phi_a} \right|$

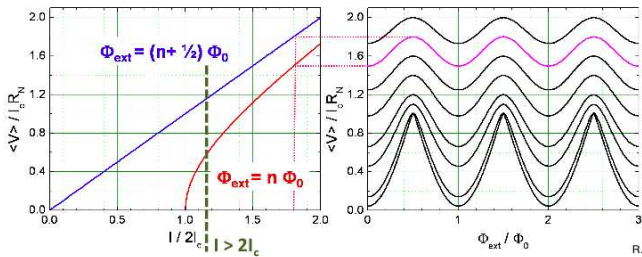
Usual operational point: $I \gtrsim 2I_c$

$$\Phi_a \simeq (2n + 1) \frac{\Phi_0}{4}$$

$$\beta \lesssim 1$$

One finds numerically

$$H \simeq \frac{I_c R_N}{\Phi_0/2} \simeq \frac{R_N}{L}$$



Figures from: R. Gross, A. Marx, Walther Meissner Institut <http://www.wmi.badw.de/teaching/Lecturenotes/>

dc SQUID: parameters

$$\langle V(t) \rangle = I_c R_N \sqrt{\left(\frac{I}{2I_c}\right)^2 - \cos^2\left(\pi \frac{\Phi_a}{\Phi_0}\right)}$$
 Flux-to-voltage converter.

Transfer function: $H = \left| \frac{\partial \langle V(t) \rangle}{\partial \Phi_a} \right|$

Usual operational point: $I \gtrsim 2I_c$

$$\Phi_a \simeq (2n + 1) \frac{\Phi_0}{4}$$

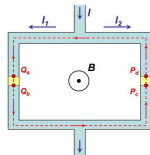
$$\beta \lesssim 1$$

One finds numerically

$$H \simeq \frac{I_c R_N}{\Phi_0/2} \simeq \frac{R_N}{L}$$

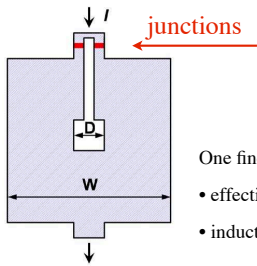
Other requirements:

- large loop (large $\Delta\Phi$ with small ΔB). But then large $L!$ \rightarrow small H and small β



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"washer" SQUIDS



“field focussing” device: the large superconducting area expels the flux, and focusses into the small hole.

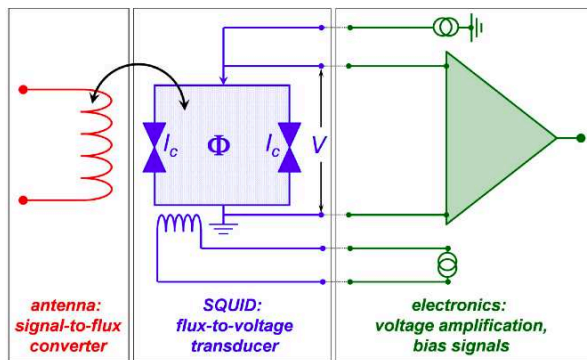
One finds:

- effective area (where field is collected) scales with W : $A_{eff} \sim D \cdot W$
- inductance scale with D : $L \approx 1.25\mu_0 D$

Caveats: W too large \rightarrow fluxons enters the superconducting area, and thermal activated motion leads to noise

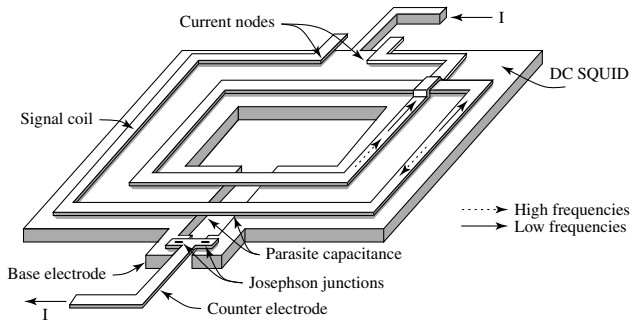
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real dc SQUID



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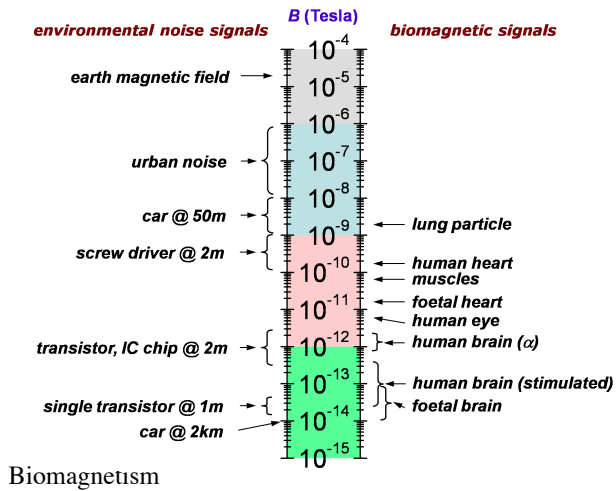
Coupling to the signal.



A spiral coil brings the signal over the SQUID area

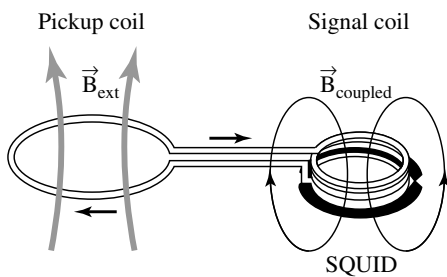
Figures from:
Fosheim-Sudbo

SQUIDs as magnetometers



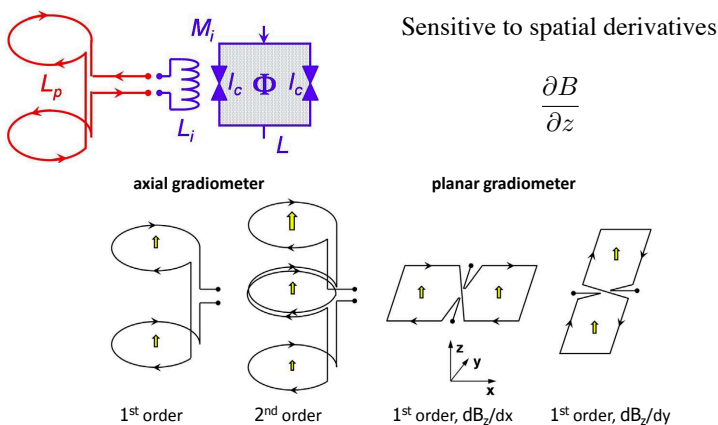
Figures from: R. Gross, A. Marx, Walther Meissner Institut <http://www.wmi.badw.de/teaching/Lecturenotes/>

Transformer configurations



Figures from: Fosheim Sudbo

Gradiometer configurations



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